

# Inverse problems for real principal type operators

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# Outline

1. Motivation
2. Results for real principal type operators
3. Methods of proof

# PDEs in a geometric setting

Consider the Laplace equation

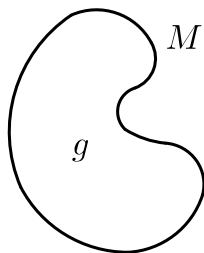
$$\Delta u = 0 \text{ in } \Omega \subset \mathbb{R}^n.$$

More generally, let  $(M, g)$  be a compact Riemannian manifold with boundary. The **Laplace-Beltrami operator**  $\Delta_g$  is

$$\Delta_g u = \sum_{j,k=1}^n \frac{1}{\sqrt{\det g}} \frac{\partial}{\partial x_j} \left( \sqrt{\det g} g^{jk} \frac{\partial u}{\partial x_k} \right)$$

where  $g = (g_{jk}(x))$ ,  $g^{-1} = (g^{jk}(x))$ .

Models an **inhomogeneous medium**.

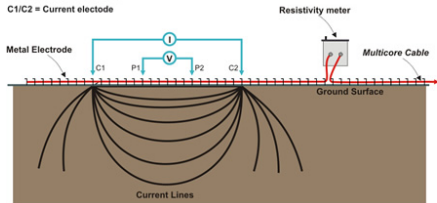


# 1. Calderón problem

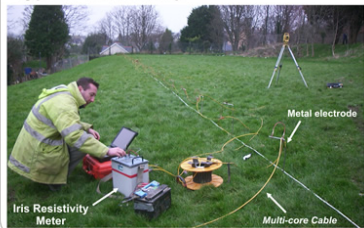
Electrical Resistivity Imaging in geophysics (1920's) [image: TerraDat]

## General resistivity principle

P1/P2 = Potential electrode  
C1/C2 = Current electrode

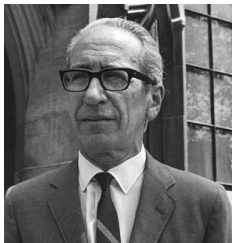


## Typical field set-up



A.P. Calderón (1980):

- ▶ mathematical formulation
- ▶ solution of the linearized problem



# 1. Calderón problem (elliptic PDE)

Laplace-Beltrami equation

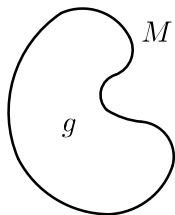
$$\begin{cases} \Delta_g u = 0 & \text{in } M, \\ u = f & \text{on } \partial M \end{cases}$$

where  $(M, g)$  is a compact Riemannian manifold with boundary ( $g \iff$  electrical conductivity).

Boundary measurements given by the **Dirichlet-to-Neumann (DN) map**

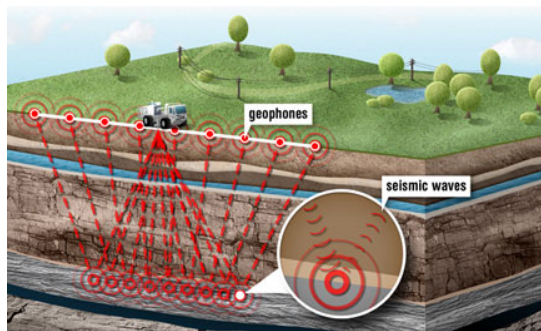
$$\Lambda_g : C^\infty(\partial M) \rightarrow C^\infty(\partial M), \quad f \mapsto \partial_\nu u|_{\partial M}.$$

**Inverse problem:** given  $\Lambda_g$ , recover  $g$ .



# 1. Gel'fand problem

Exploration seismology:



I.M. Gel'fand (1950s):

- ▶ 1D inverse scattering (Gel'fand, Levitan, Marchenko)
- ▶ proposed the question for  $n \geq 2$

## 2. Gel'fand problem (hyperbolic PDE)

Wave equation ( $g \leftrightarrow$  sound speed)

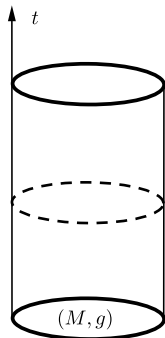
$$\begin{cases} (\partial_t^2 - \Delta_g)u = 0 & \text{in } M \times (0, T), \\ u = f & \text{on } \partial M \times (0, T) \\ u|_{\{t < 0\}} = 0 \end{cases}$$

where  $(M, g)$  is compact with boundary.

Boundary measurements given by the  
**hyperbolic DN map**

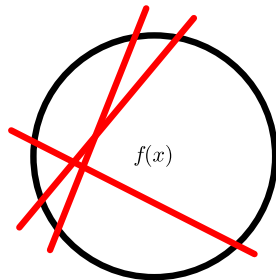
$$\Lambda_g^{\text{Hyp}} : f \mapsto \partial_\nu u|_{\partial M \times (0, T)}.$$

**Inverse problem:** given  $\Lambda_g^{\text{Hyp}}$ , recover  $g$ .



### 3. X-ray transform problems

X-ray computed tomography:



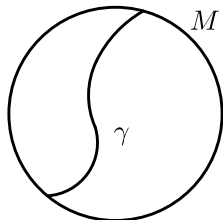
Recover a function  $f(x)$  from its integrals over straight lines [Radon 1917]. Seismic imaging: lines are replaced by geodesics.



### 3. X-ray transform / scattering relation

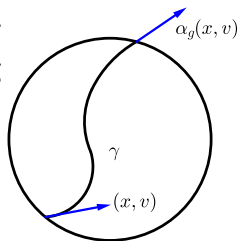
Try to recover a function  $f$  in  $(M, g)$  from its *geodesic X-ray transform*  $I f$ , where

$$I f(\gamma) = \int_{\gamma} f dt, \quad \gamma \text{ maximal geodesic.}$$



Related **scattering rigidity** problem: recover  $(M, g)$  from its **scattering relation**  $\alpha_g$ , relating initial and final data of maximal geodesics:

$$\alpha_g : (x, v) \mapsto \alpha_g(x, v)$$



Can formulate both questions in terms of (transport) PDEs. These are **highly nonlinear questions** related to **linear PDEs**!

# Connections

Unexpected connections in special geometries:

- ▶ Calderón problem reduces to geodesic X-ray transform  
[Dos Santos-Kenig-S-Uhlmann 2009]
- ▶ Calderón problem reduces to Gel'fand problem  
[Dos Santos-Kurylev-Lassas-S 2016]
- ▶ scattering rigidity problem reduces to Calderón problem  
[Pestov-Uhlmann 2005]

What are the general structural conditions and mechanisms behind this?

# Goal

Propose to study inverse problems for **general** differential operators. Hope to understand:

- ▶ structural conditions for treating classes of operators
- ▶ fundamental mechanisms for solving inverse problems
- ▶ the extent to which it is possible to push existing methods.

Approach in the spirit of general theory for linear PDEs [Hörmander 1983–1985]. Earlier results for constant coefficients [Isakov 1991, ...].

We will consider **variable coefficients** (=microlocal analysis).

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# Real principal type operators

Let  $M$  be compact with boundary. A differential operator  $P$  on  $M$ , of order  $m \geq 1$ , is **real principal type** if

- ▶ it has **real principal symbol**  $\sigma_{\text{pr}}(P) = p_m$ ,<sup>1</sup> and
- ▶ the null bicharacteristic flow is **nontrapping**.

**Null bicharacteristic curves**  $\gamma(t) = (x(t), \xi(t))$  are integral curves of Hamilton field  $H_{p_m}$  in  $p_m^{-1}(0)$ . They solve the ODE

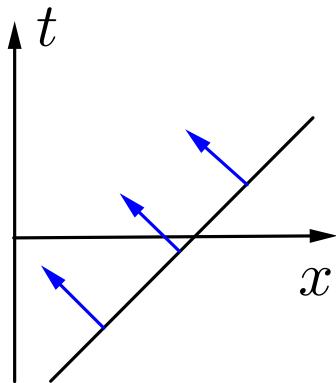
$$\begin{cases} \dot{x}(t) &= \nabla_{\xi} p_m(x(t), \xi(t)), \\ \dot{\xi}(t) &= -\nabla_x p_m(x(t), \xi(t)). \end{cases}$$

**Nontrapping** means that any such  $\gamma(t)$  reaches  $\partial M$  in finite time in both directions.

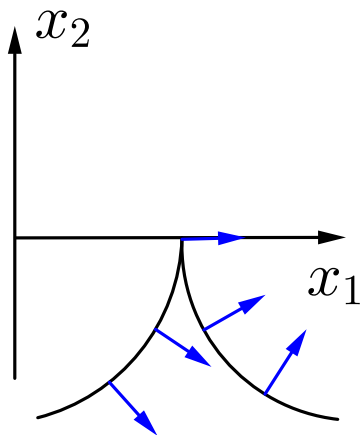
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<sup>1</sup>If  $P = \sum_{|\alpha| \leq m} p_{\alpha}(x) D^{\alpha}$ , then  $\sigma_{\text{pr}}(P) = \sum_{|\alpha|=m} p_{\alpha}(x) \xi^{\alpha}$ .

## Null bicharacteristics



Wave operator  $\partial_t^2 - \Delta$



Tricomi operator  $x_2 D_{x_1}^2 + D_{x_2}^2$

# Real principal type operators

Examples:

- ▶ real vector fields with no trapped integral curves
- ▶ wave operator in  $M \times (0, T)$ , Lorentzian wave operators with nontrapping condition, strictly hyperbolic operators
- ▶ Tricomi type operators, e.g.  $x_2 D_{x_1}^2 + D_{x_2}^2$
- ▶ Schrödinger operator  $i\partial_t + \Delta$ , plate equation  $\partial_t^2 + \Delta^2$  with suitable (anisotropic) weighting for  $\partial_t$

Real principal type operators can be microlocally conjugated to normal form  $D_{x_1}$ . Singularities of solutions propagate along null bicharacteristics, solvability theory for  $Pu = f$

[Duistermaat-Hörmander 1972].

## Boundary measurements

It is not clear how to define an analogue of DN map for a general operator. However, we consider the **Cauchy data set**

$$C_P = \{(u|_{\partial M}, \dots, \nabla^{m-1} u|_{\partial M}); Pu = 0 \text{ in } M, u \in H^m(M)\}.$$

This is equivalent to knowing the DN map e.g. in the Calderón and Gel'fand problems.

**Inverse problem:** given  $C_P$ , determine information about  $P$ .

From now on, all operators will be **real principal type** in  $M$ .

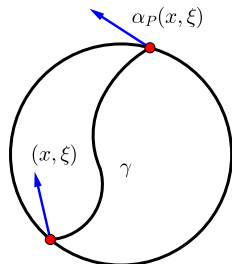


# Determining (sub)principal information

## Theorem 1

If  $C_{P_1} = C_{P_2}$  and if  $P_1 = P_2$  to infinite order on  $\partial M$ , then

$$\alpha_{P_1} = \alpha_{P_2}$$



where  $\alpha_P$  is the **bicharacteristic scattering relation**, mapping an initial point of a **maximal null bicharacteristic** to its final point.

Moreover, if  $P_1$  and  $P_2$  have the same principal symbol, then

$$\exp \left[ i \int \sigma_{\text{sub}}(P_1)(\gamma(t)) dt \right] = \exp \left[ i \int \sigma_{\text{sub}}(P_2)(\gamma(t)) dt \right]$$

for any **maximal null bicharacteristic**  $\gamma$  in  $T^*M$ .

## Lower order coefficients

The conclusion  $\exp[i \int \dots] = \exp[i \int \dots]$  is equivalent with

$$\int \sigma_{\text{sub}}(P_1)(\gamma(t)) dt = \int \sigma_{\text{sub}}(P_2)(\gamma(t)) dt \pmod{2\pi\mathbb{Z}}.$$

This is related to the **Aharonov-Bohm effect** in determining subprincipal terms on domains with nontrivial topology. For lower order coefficients, this effect does not appear:

### Theorem 2 (Bicharacteristic ray transforms)

If  $C_{P+Q_1} = C_{P+Q_2}$  where  $Q_j$  are operators of order  $\leq m - 2$ , then

$$\int \sigma_{\text{pr}}(Q_1)(\gamma(t)) dt = \int \sigma_{\text{pr}}(Q_2)(\gamma(t)) dt$$

for any maximal null bicharacteristic  $\gamma$  in  $T^*M$ .

# Real principal type operators

The results are general: they extend results for wave equations [Rakesh-Symes 1988, . . . , Stefanov-Yang 2018], and are valid for

- ▶ operators of any order, with **real principal symbol** and **nontrapping** condition (no wellposedness assumptions)
- ▶ **any maximal bicharacteristic**, even with cusps (Tricomi) and tangential reflections

However, the results are **conditional**: in order to recover coefficients of  $P$ , one still needs to analyze the scattering relation  $\alpha_P$  or bicharacteristic ray transforms.

# Boundary determination

Determine Taylor series of coefficients of  $P$  at null points  $(x, \xi) \in T^*(\partial M)$ , based on zeros of **characteristic polynomial**

$$t \mapsto p_m(x, \xi + t\nu).$$

Two methods:

1. **Elliptic region.** If there is a simple non-real zero, use exponentially decaying solutions (analogue of boundary determination for Laplace equation).
2. **Hyperbolic region.** If there are two distinct real zeros, use solutions concentrating near two null bicharacteristics (analogue of boundary determination for wave equation).

# Boundary determination

## Theorem 3 (Determining Taylor series of a potential)

If  $V_1, V_2 \in C^\infty(M)$  and  $C_{P+V_1} = C_{P+V_2}$ , then

$$\nabla^k V_1(x_0) = \nabla^k V_2(x_0), \quad k \geq 0,$$

at any  $x_0 \in \partial M$  so that for some  $\xi \in T_x^*(\partial M)$ , the map  $t \mapsto p_m(x_0, \xi + t\nu)$  either has a simple non-real root, or two distinct real roots<sup>1</sup>.

In particular, if  $M$  and  $V_j$  are **real-analytic** and there is one such  $x_0$ , then  $V_1 = V_2$  **everywhere in  $M$** .

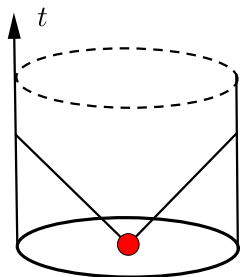
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<sup>1</sup>with corresponding bicharacteristics intersecting nicely at  $x_0$

# Boundary determination

Observations:

- ▶ boundary determination in general not possible for  $m = 1$
- ▶ even for wave equation, can do boundary determination in the elliptic region as for elliptic operators (local argument)
- ▶ boundary determination in the hyperbolic region is global in character



$$P = \partial_t^2 - \Delta + V$$

# Nonlinear equations

If  $q \in C^\infty(M)$ , consider the semilinear equation

$$Pu + q(x)u^k = 0 \text{ in } M.$$

Let  $C_q^{\text{small}}$  be the Cauchy data set for small solutions.

## Theorem 4 (Semilinear equations)

Let  $q_1, q_2 \in C^\infty(M)$  and  $k \geq 3$ . If  $C_{q_1}^{\text{small}} = C_{q_2}^{\text{small}}$ , then  $q_1 = q_2$  in  $B$  where

$B = \{x \in M; \text{ there are two null bicharacteristics that intersect only once at } x \text{ transversally}\}.$

Nonlinearity helps (proof fails if  $k = 1$ )! Wave equations:

[Kurylev-Lassas-Uhlmann 2018, Lassas-Uhlmann-Wang 2018, Hintz-Uhlmann 2018]

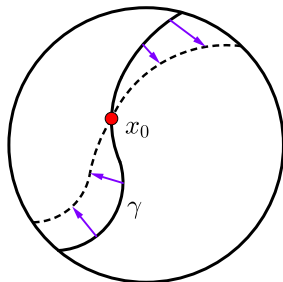
# Nonlinear equations

Can recover the coefficient  $q(x)$  in the set

$B = \{x_0 \in M; \text{there are two null bicharacteristics that intersect only once at } x_0 \text{ transversally}\}.$

If there is a nice<sup>1</sup> bicharacteristic  $\gamma(t)$  through  $x_0$  having a variation field only vanishing at  $t = 0$ , can recover  $q(x_0)$ .

Works e.g. if some  $\gamma(t)$  through  $x_0$  has "no conjugate points". May fail if there is a "maximally conjugate" point.



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<sup>1</sup>nontangential, no cusp at  $x_0$ ,  $x(t)$  does not self-intersect



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# Methods

1. Use Cauchy data of **special solutions** concentrating along a **null bicharacteristic** (propagation of singularities).
2. For boundary determination, also use **exponentially decaying solutions** concentrating at a boundary point.
3. Use **integral identities** and a **mix-and-match construction** to pass from Cauchy data set  $C_P$  to scattering relation / bicharacteristic ray transforms / pointwise information.

# Quasimode construction

## Theorem 5

Let  $P$  have **real principal symbol** in an open mfld  $X$ , and let  $\gamma : [0, T] \rightarrow T^*X$  be an **injective** null bicharacteristic segment. There is  $u = u_h \in C_c^\infty(X)$  with

$$WF_{\text{scl}}(u) = \gamma([0, T]), \quad WF_{\text{scl}}(Pu) = \gamma(0) \cup \gamma(T),$$

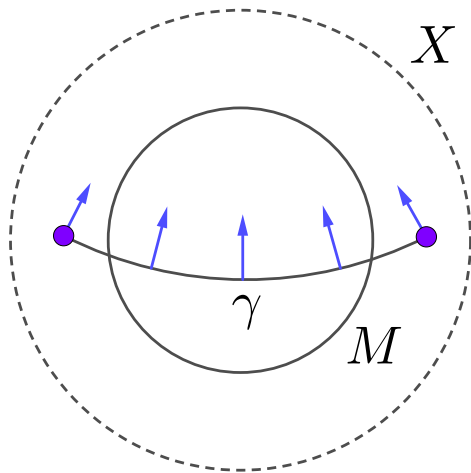
having semiclassical defect measure (with  $c_\gamma(t) > 0$ )

$$\lim_{h \rightarrow 0} (Op_h(a)u_h, u_h)_{L^2(X)} = \int_0^T a(\gamma(t))c_\gamma(t) dt$$

whenever  $a \in S^0$  vanishes near endpoints of  $\gamma$ .

Semiclassical counterpart of [\[Duistermaat-Hörmander 1972\]](#).

## Quasimode construction



$$WF_{\text{scl}}(Pu) = \gamma(0) \cup \gamma(T) \implies Pu = O(h^\infty) \text{ in } M$$

# Methods for constructing quasimodes

1. Locally, enough to use **geometrical optics**:

$$u_h(x) = e^{i\varphi(x)/h} a(x)$$

where  $\varphi$  is real and  $p_m(x, d\varphi(x)) = 0$  (**eikonal equation**).

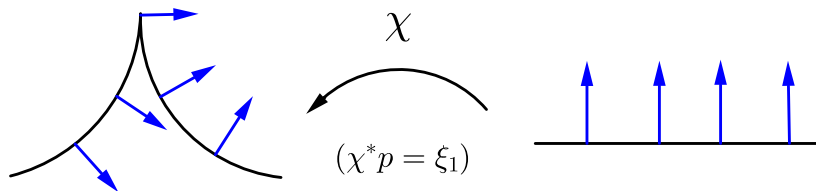
2. No cusps ( $\dot{x}(t) \neq 0$ ): can use a **Gaussian beam** construction

$$u_h(x) = e^{i\Phi(x)/h} a(x)$$

where  $\Phi$  is **complex** and solves eikonal equation to infinite order on the curve  $x(t)$ . Main point:  $\nabla^2\Phi|_{x(t)}$  solves a matrix Riccati equation and  $\text{Im}(\nabla^2\Phi) > 0$ .

## Quasimodes [Duistermaat-Hörmander 1972]

3. If  $\gamma(t)$  is injective but may have cusps, can straighten  $\gamma(t)$  in phase space by a canonical transformation  $\chi$ .



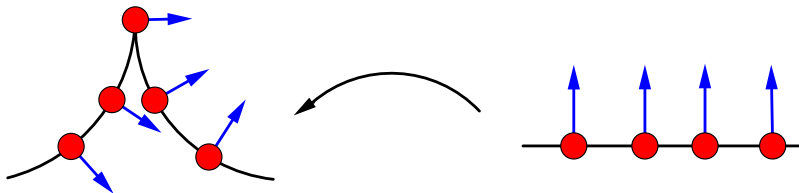
Multiply  $P$  by an elliptic  $\Psi$ DO so that  $P$  becomes of order 1. Construct Fourier integral operators  $A, B$  such that

$$BPA \approx D_{x_1} \text{ microlocally near } \gamma.$$

Quasimode  $U$  for  $D_{x_1} \implies u = AU$  is a quasimode for  $P$ .

## Quasimodes (direct construction)

4. Think of quasimodes as superpositions of wave packets  $\approx e^{i\frac{\xi(t)\cdot(x-x(t))}{h}} e^{-\frac{|x-x(t)|^2}{2h}}$  at  $x(t)$  oscillating in direction  $\xi(t)$ .



Look for  $u_h$  with  $Pu_h = O(h^\infty)$  directly in the form

$$u_h(x) = \int_0^T e^{i\Phi(x,t)/h} a(x,t) dt.$$

Cf. Gaussian beam construction along  $(x(t), t)$  in  $X \times \mathbb{R}$ .

# Future directions

1. Inversion of scattering relation  $\alpha_P$ ? If  $P = X_g$ ,<sup>1</sup> studied in [Pestov-Uhlmann 2005, Stefanov-Uhlmann-Vasy 2017].
2. Inversion of bicharacteristic ray transform? Cf. geodesic ray transform [Uhlmann-Vasy 2016, Paternain-S-Uhlmann 2015] ( $P = X_g$ ), and light ray transform [Lassas et al 2019] ( $P = \square_g$ ).
3. Results for mild trapping? For  $P = X_g$  and hyperbolic trapping, studied in [Guillarmou / Guillarmou-Monard 2017].
4. Can one associate a symbol directly to  $C_P$ ?
5. The results are in the spirit of using singularities of the integral kernel of DN map. Can one extract information from the  $C^\infty$  part of the kernel?

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<sup>1</sup>geodesic vector field on unit sphere bundle