

Abelian and non-Abelian X-ray transforms: mapping properties and Bayesian inversion

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UNIVERSITY OF CALIFORNIA
SANTA CRUZ



- 1 Introduction
- 2 Main results
- 3 A numerical illustration
- 4 Elements of proof: mapping properties of unattenuated X-ray
 - The 'classical' functional setting $L^2(M) \rightarrow L^2_{\mu}(\partial_+ SM)$
 - The Euclidean disk: setting $L^2(\mathbb{D}) \rightarrow L^2(\partial_+ S\mathbb{D})$
- 5 Elements of proof: mapping properties of attenuated X-ray

Outline

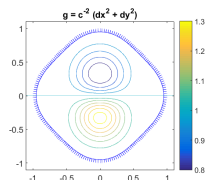
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The geodesic X-ray transform

(M, g) , ∂M strictly convex.

$\partial_+ SM$: “inward” boundary (‘fan-beam’).

Geodesics: $\gamma_{x,v}(t)$.



Geodesic X-ray transform of f :

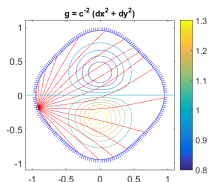
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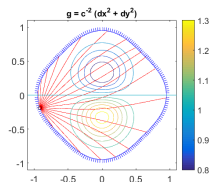
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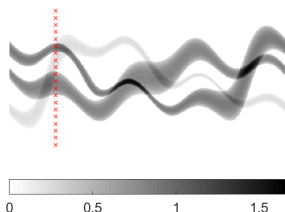
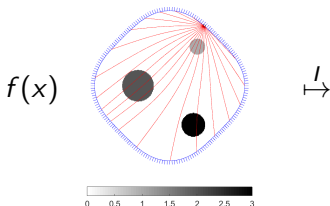
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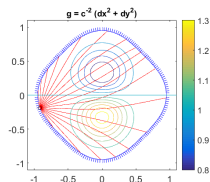


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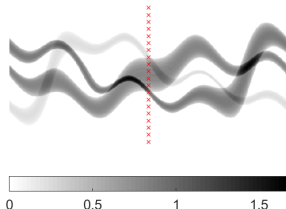
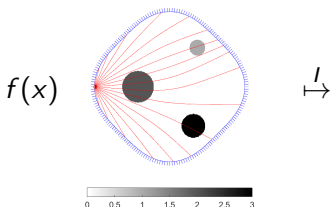
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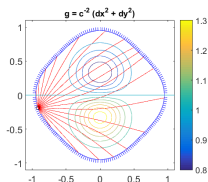


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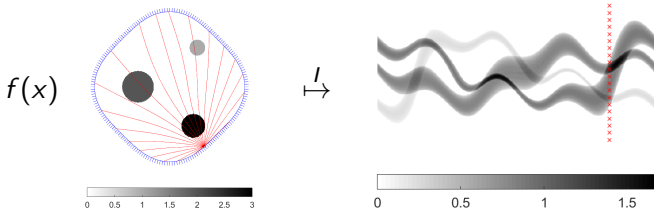
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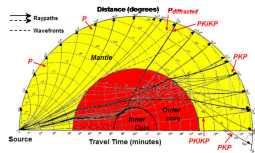
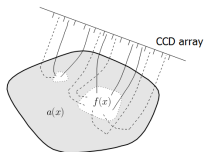
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Goal: recover f from If .

Applications of geodesic X-ray transform

- 1 Radon transform and X-ray CT
- 2 SPECT and tomography in media with variable index of refraction
- 3 Seismology and travel-time tomography.

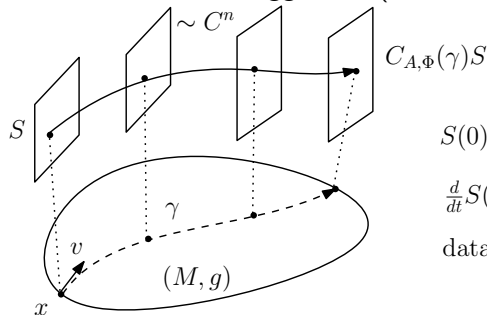


The non-abelian X-ray transform

$E = M \times \mathbb{C}^n$ trivial bundle over (M, g) a simple Riem. surface

A : hermitian connection on E (= matrix of one-forms)

Φ : skew-hermitian Higgs field ($\Phi : M \rightarrow \mathfrak{u}(n)$)



$$S(0) = S$$

$$\frac{d}{dt}S(t) + (A_{\gamma(t)}(\dot{\gamma}(t)) + \Phi_{\gamma(t)})S(t) = 0$$

$$\text{data: } S(\ell(\gamma)) = C_{A, \Phi}(\gamma)S$$

Inverse problem:

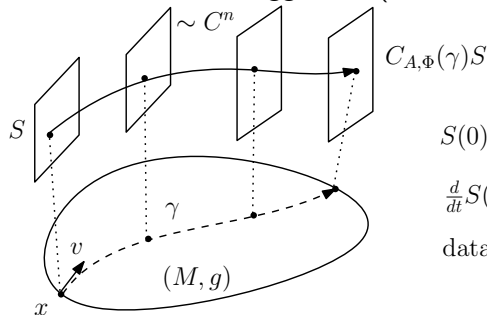
- to recover (A, Φ) from $C_{A, \Phi}$ modulo natural obstruction.
- if $A \equiv 0$: to recover Φ from C_{Φ} .

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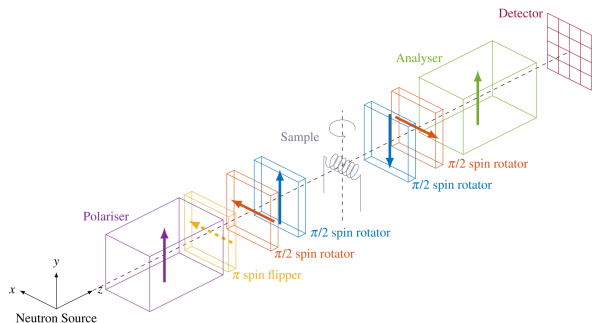
Problem (Polarimetric Neutron Tomography)

To recover a magnetic field from neutron spin-in to spin-out map (“scattering data” of the magnetic field).

While a neutron's trajectory $\gamma(t)$ is unaffected by a magnetic field B , its **spin** S evolves according to

$$\dot{S}(t) = B(t) \times S(t).$$

$B \leftrightarrow \mathfrak{so}(3)$ -valued Higgs field.



Src: Sales et al, “3D Polarimetric Neutron Tomography of Magnetic Fields”, 2017.

Bayesian approach to noisy inversion 1/2

We consider the problem of recovering a matrix field Φ_0 from

$$Y_j = C_{\Phi_0}(\gamma_j) + \varepsilon_j, \quad \varepsilon_j \sim \mathcal{N}(0, \sigma^2), \quad 1 \leq j \leq N.$$

where $\{\gamma_j\}_{j=1}^N$ are chosen uniformly at random in fan-beam coordinates. Denote $D_N = \{(\gamma_j, Y_j)\}_{1 \leq j \leq N}$ the data set.

Given a **prior** model for Φ and a **noise model**, Bayes' formula gives the proba. density of the posterior random variable $\Phi|D_N$:

$$P(\Phi|D_N) = \frac{P(D_N|\Phi)P(\Phi)}{P(D_N)} \propto \underbrace{P(D_N|\Phi)}_{\text{likelihood}} \cdot \underbrace{P(\Phi)}_{\text{prior}}$$

With a Gaussian prior and Gaussian noise model, one may arrive at the log-posterior distribution:

$$-\log \Pi(\Phi|D_N) = \frac{1}{\sigma^2} \sum_{j=1}^N |Y_j - C_{\Phi}(\gamma_j)|_F^2 + \|\Phi\|_{H^\alpha}^2 + C.$$

Note: Not gaussian since $\Phi \mapsto C_{\Phi}$ is non-linear.

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Bayesian approach to noisy inversion 2/2

Rather than looking at the entire posterior object $\Phi|D_N$, consider ONE of its moments $\langle \Phi, \psi \rangle | D_N$, where ψ is a 'test' field.

Important questions:

- 1 What relevant estimators $\hat{\Psi}_N$ of $\langle \Phi_0, \psi \rangle$ can we extract from this posterior distribution ? Examples: mean and MAP, when they exist.
- 2 Are they easily computable ?
- 3 Consistency: does $\hat{\Psi}_N \rightarrow \langle \Phi_0, \psi \rangle$ as $N \rightarrow \infty$ in any sense ?
- 4 UQ: does the posterior density "contract" about $\hat{\Psi}_N$ so one can get confidence/credible sets ?

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Main messages of this talk

- 1 Making UQ statements for inverse problems requires a refined understanding of mapping properties of the forward and/or the 'normal' (linearized) operator.
- 2 For X-ray problems on manifolds with boundary, these mapping properties (the natural Hilbert scales involved and their Fréchet limits) are highly sensitive to the choice of certain weights.
- 3 Said choice of weights can come from: theoretical tractability ("choose the easiest one !") or noise model ("choose the practical one !").

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Main results 1/3

Theorem (M', 2019, arXiv preprint 1910.13691)

On M the Euclidean disk, if $I_0: L^2(M) \rightarrow L^2(\partial_+ SM)$ is the X-ray transform with adjoint I_0^* , we have

$$I_0^* I_0: C^\infty(M) \xrightarrow{\cong} C^\infty(M), \quad I_0^* I_0: \tilde{H}^s(M) \xrightarrow{\cong} \tilde{H}^{s+1}(M), \quad s \geq 0,$$

where $\tilde{H}^s := D(\mathcal{L}^s)$ upon defining

$$\mathcal{L} = -((1 - \rho^2)\partial_\rho^2 + (\rho^{-1} - 3\rho)\partial_\rho + \rho^{-2}\partial_\omega^2) + 1.$$

- The L^2 topology on $\partial_+ SM$ is not the usual, 'symplectic' one.
- Generalizations to special cases of simple surfaces also exist.

Main results 2/3

Theorem (M'-Nickl-Paternain, 2020, arXiv preprint 2007.15892)

On M the Euclidean disk, let $\Theta \in C_c^\infty(M, \mathfrak{so}(n))$, and let the attenuated X-ray transform

$$I_\Theta : L^2(M, \mathbb{C}^n) \rightarrow L^2(\partial_+ SM, \mathbb{C}^n).$$

Then $I_\Theta^* I_\Theta$ is an isomorphism

$$C^\infty(M, \mathbb{C}^n) \xrightarrow{\cong} C^\infty(M, \mathbb{C}^n), \quad \tilde{H}^s(M, \mathbb{C}^n) \xrightarrow{\cong} \tilde{H}^{s+1}(M, \mathbb{C}^n), \quad s \geq 0.$$

- based on setting up a suitable Fredholm setting, using the \tilde{H}^s scale and $\Theta \equiv 0$ as reference case.

Main result 3/3 - application to UQ

Consider the problem of recovering $\Phi_0 \in C_c^\infty(M, \mathfrak{so}(n))$ from

$$Y_j = C_{\Phi_0}(\gamma_j) + \varepsilon_j, \quad \varepsilon_j \sim \mathcal{N}(0, \sigma^2), \quad 1 \leq j \leq N.$$

As above, let $D_N := \{\gamma_j, Y_j\}_{j=1}^N$, and choose a prior among a flexible class and a test field $\psi \in C^\infty(M, \mathfrak{so}(n))$. Denote $\hat{\Phi}_N$ the posterior mean (from [M.-Nickl-Paternain, CPAM '20], it exists and converges to Φ_0 as $N \rightarrow \infty$).

Theorem (M.-Nickl-Paternain, 2020, arXiv preprint 2007.15892)

We have as $N \rightarrow \infty$ and in $P_{\Phi_0}^N$ -probability, the weak convergence

$$\sqrt{N} \langle \Phi - \hat{\Phi}_N, \psi \rangle | D_N \rightarrow^d \mathcal{N}(0, \|I_{\Theta_0}(I_{\Theta_0}^* I_{\Theta_0})^{-1} \psi\|_{L^2(\partial_+ SM)}^2).$$

- I_{Θ_0} is an attenuated X-ray transform where Θ_0 depends on Φ_0 , related to the linearization of the map $\Phi \mapsto C_\Phi$ at Φ_0 .

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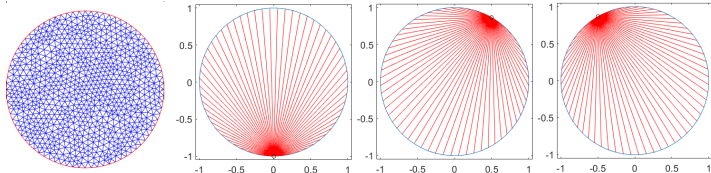
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Computational domain, unknown

M : Euclidean unit disk.

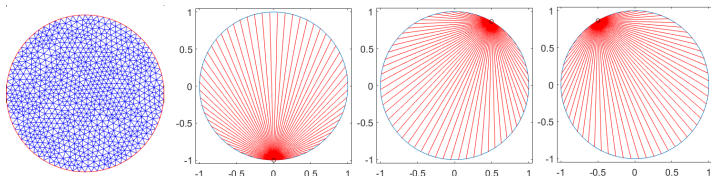
Unstructured mesh with 886 nodes.



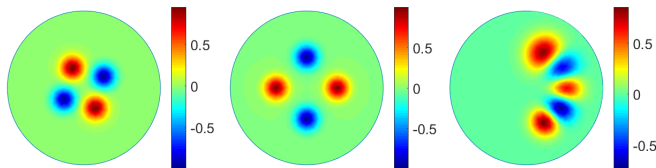
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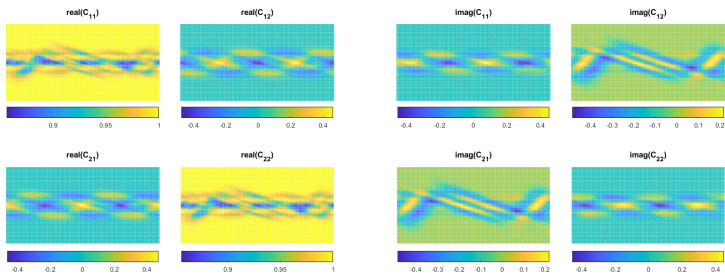
Magnetic field (3 components):



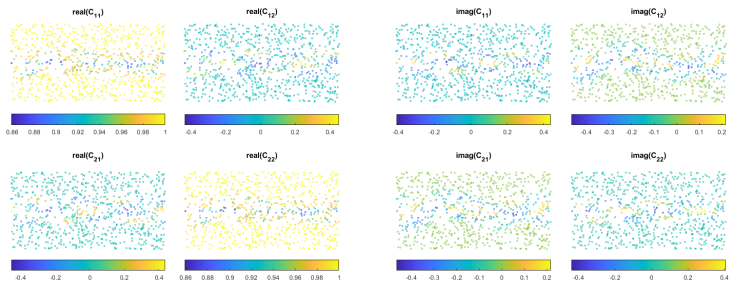
$$\text{(made } \mathfrak{su}(2)\text{-valued via } \mathbb{R}^3 \ni (a, b, c) \mapsto \frac{1}{2} \begin{bmatrix} ia & b + ic \\ -b + ic & -ia \end{bmatrix} \text{)}$$

Forward data + sampling

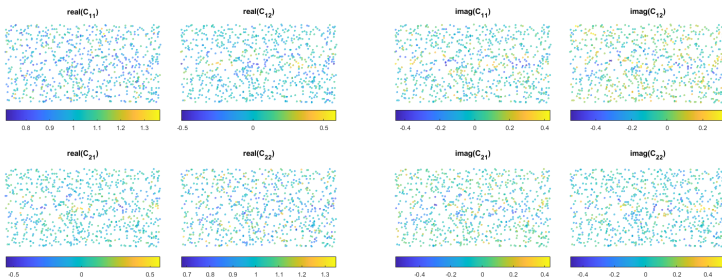
Noiseless



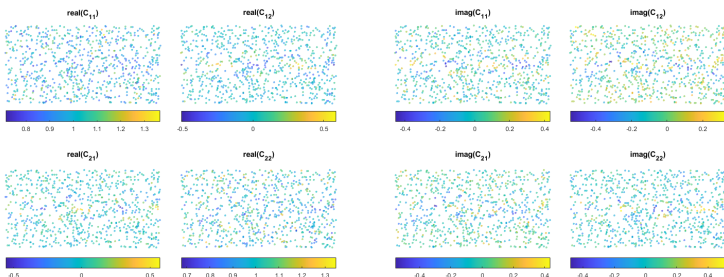
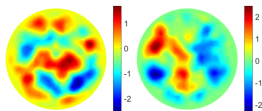
Forward data + sampling

Noiseless, randomly sampled ($N = 800$)

Forward data + sampling

Noisy ($\sigma = 0.1$), randomly sampled ($N = 800$)

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Noisy ($\sigma = 0.1$), randomly sampled ($N = 800$)Prior (Matérn) parameters: $\nu = 3$, $\ell = 0.2$. Samples:

Computation of the posterior mean by MCMC

We compute a family of posterior draws using **preconditioned Crank-Nicolson** [Cotter, Stuart, Roberts & White '13].

Fix $\delta \in (0, 1/2)$, $\Phi_0 = 0$, and for $n = 0 : N_s$ do:

- 1 Draw $\Psi \sim \Pi$ and set $p_n = \sqrt{1 - 2\delta} \Phi_n + \sqrt{2\delta} \Psi$.
- 2 With $\ell(\Phi) = \frac{1}{\sigma^2} \sum_{j=1}^N |Y - C_\Phi(\gamma_j)|_F^2$ the log-likelihood,

$$\text{set } \Phi_{n+1} = \begin{cases} p_n & \text{with proba. } 1 \wedge \exp(\ell(\Phi_n) - \ell(p_n)), \\ \Phi_n & \text{otherwise.} \end{cases}$$

- ⊗ Visualize $\hat{\Phi} = \frac{1}{N_s} \sum_{n=1}^{N_s} \Phi_n$ and histograms of moments.

• One can show that $\{\Phi_n\}_n$ forms a Markov chain with unique invariant measure $\Pi(\cdot | (Y, x)_{n=1}^N)$.

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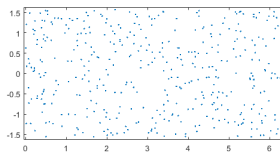
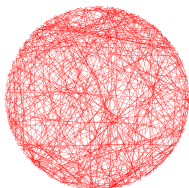
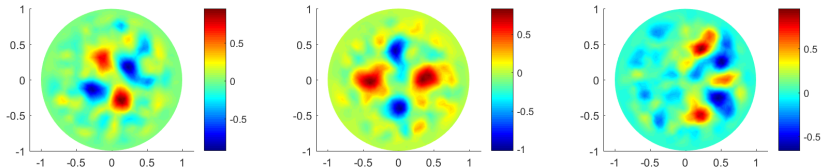
$$\text{set } \Phi_{n+1} = \begin{cases} p_n & \text{with proba. } 1 \wedge \exp(\ell(\Phi_n) - \ell(p_n)), \\ \Phi_n & \text{otherwise.} \end{cases}$$

- ⊗ Visualize $\hat{\Phi} = \frac{1}{N_s} \sum_{n=1}^{N_s} \Phi_n$ and histograms of moments.

- One can show that $\{\Phi_n\}_n$ forms a Markov chain with unique invariant measure $\Pi(\cdot | (Y_i, x_i)_{i=1}^N)$.
- No inversion required, only forward solves !
- [Hairer-Stuart-Vollmer '14] prove non-asymptotic mixing $Law(\Phi_n) \rightarrow \Pi(\cdot | Y)$ under conditions on $\ell(\Phi)$ that do not require convexity.

Illustration of consistency - posterior mean

$$N = 400, \sigma = 0.05, \delta = 2.5 \cdot 10^{-5}$$

MCMC sample average over 10^5 iterations

truth:

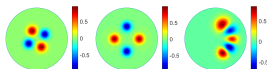
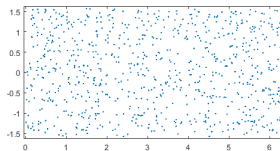
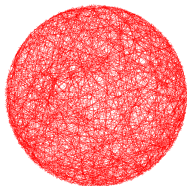
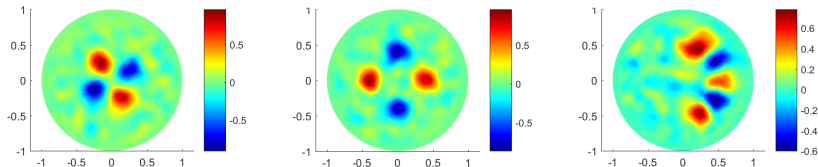


Illustration of consistency - posterior mean

$$N = 800, \sigma = 0.05, \delta = 2.5 \cdot 10^{-5}$$

MCMC sample average over 10^5 iterations

truth:

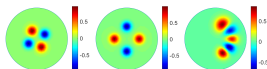
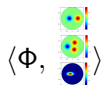
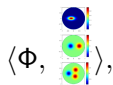
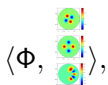
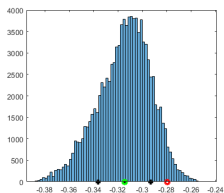
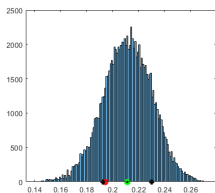
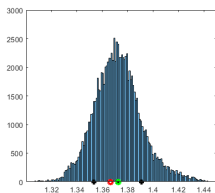


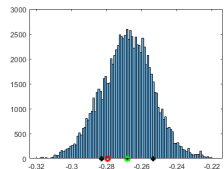
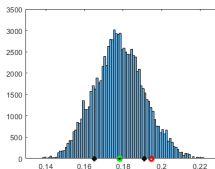
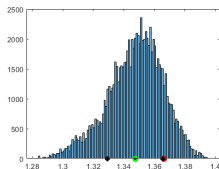
Illustration of approximate normality - histograms



$N = 600, 10^6$ iterations



$N = 1000, 10^6$ iterations



Outline

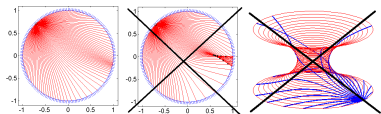
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The case of simple surfaces

Simple = ∂M strictly convex + no conjugate points + no geodesic of infinite length.



Recovery of f from If is injective [Mukh. '75], ill-posed of order 1/2 in that one may derive the stability estimate

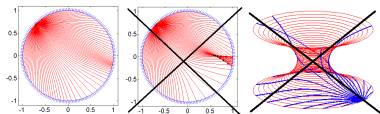
$$\|f\|_{L^2(M)} \leq C \|I^\# If\|_{H^1(\tilde{M})}, \quad (I^\# : L^2(M) \rightarrow L^2_\mu(\partial_+ SM) \text{ adjoint})$$

where \tilde{M} is a simple extension of M . [Stef.-Uhl. '04]

Question: Can one obtain mapping properties of $I^\#$ that do not require an extension ?

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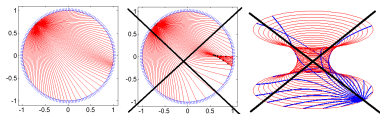
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Extendibility property

In the $L^2(M) \rightarrow L^2_\mu(\partial_+ SM)$ setting, the normal operator looks like

$$I^\# I f(x) = 2 \int_{S_x} \int_0^{\tau(x,v)} f(\gamma_{x,v}(t)) dt dS(v).$$

If \tilde{M} is a simple extension of M , one could define $\tilde{I}^\# I$ similarly, and notice that

$$r_M \circ \tilde{I}^\# I \circ e_M = I^\# I.$$

Moreover, $\tilde{I}^\# I \in \Psi_{ell}^{-1}(\tilde{M})$ and satisfies a $-1/2$ transmission condition at ∂M , a symmetry condition on its full symbol expansion relating $\sigma(x, \nu_x)$ and $\sigma(x, -\nu_x)$ at every point $x \in \partial M$.

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Mapping properties of I^*I

Theorem (M.-Nickl-Paternain, AoS '19)

The map $I^\sharp I$ is an isomorphism in the settings below:

$$(i) \ I^\sharp I: d^{-1/2} C^\infty(M) \rightarrow C^\infty(M), \quad d(x) = \text{dist}(x, \partial M)$$

$$(ii) \ I^\sharp I: H^{-1/2(s)}(M) \rightarrow H^{s+1}(M), \quad s > -1, \quad (\text{bi-continuous}).$$

$H^{\mu(s)}(M)$: Hörmander μ -transmission spaces.

$$\cap_s H^{\mu(s)}(M) = d^\mu C^\infty(M).$$

Proof of (i) (sketch):

- 1 $I^\sharp I$ is Fredholm. Uses the μ -transmission property for Ψ DOs.
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Comments

- ⊕ It's sharp and does not require extension.
- ⊕ Classical Sobolev scales cannot be used everywhere.
- ⊖ The Hörmander transmission spaces aren't great to work with (not clear what $I(H^{-1/2(s)}(M))$ looks like)
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- Can we get $C^\infty(M) \rightarrow C^\infty(M)$ isomorphism ?
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Hunch: change the weight on the co-domain because in fact, $I: L^2(M) \rightarrow L^2(\partial_+ SM)$ is bounded. It is also **the** functional setting where the SVD is known in the Euclidean disk !

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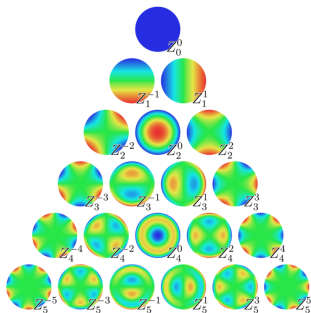
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The Euclidean disk

The SVD has been long known [Cormack, Maass, Louis...].

Zernike polynomials:

$$Z_{n,k}, n \in \mathbb{N}_0, 0 \leq k \leq n.$$



Uniquely defined through:

- $Z_{n,0} = z^n.$
- $\partial_{\bar{z}} Z_{n,k} = -\partial_z Z_{n,k-1},$
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- $Z_{n,k}|_{\partial M}(e^{i\beta}) = e^{i(n-2k)\beta}.$

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[Kazantzev-Bukhgeym '07]

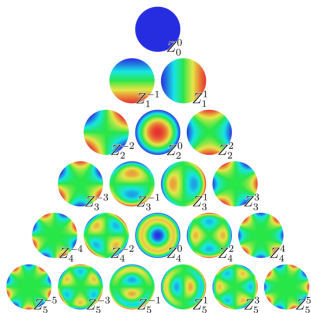
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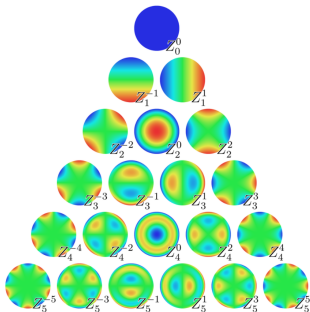
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Denote I^* the adjoint in this setting ($I^* = I^\# \frac{1}{\mu}$). Let $T = \partial_\beta - \partial_\alpha$ and $\mathcal{L} := -((1 - \rho^2)\partial_\rho^2 + (1/\rho - 3\rho)\partial_\rho + 1/\rho^2\partial_\omega^2) + 1$. Facts:

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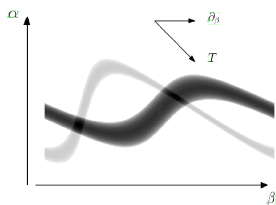
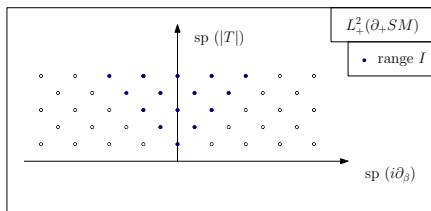
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- $\tilde{H}^1(\mathbb{D}) \not\supseteq H^1(\mathbb{D})$.

Euclidean disk: on the data side [M. '19]

The relation $I \circ \mathcal{L} = (-T^2) \circ I$ indicates that smoothness in \mathcal{L} translates into smoothness along $(-T^2)$.

Define $H_{T,+}^{1/2}(\partial_+ S\mathbb{M}) = \{w \in L_+^2, (-T^2)^{k/2} w \in L_+^2\}$ to obtain

$$\|If\|_{H_{T,+}^{k+1/2}(\partial_+ S\mathbb{D})} = c \|f\|_{\tilde{H}^k(\mathbb{D})}, \quad \forall f, \forall k.$$



Similar anisotropic scales constructed in [Natterer, Assylbekov-Stefanov '19, Paternain-Salo '19]

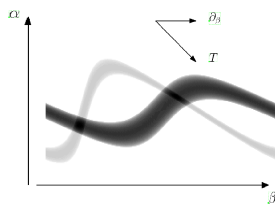
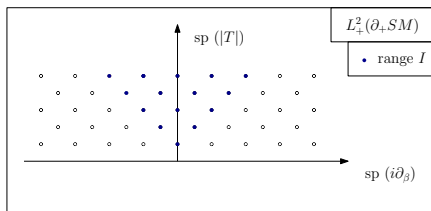
There also exists a projection operator onto $\text{range}(I)$.

Euclidean disk: on the data side [M. '19]

The relation $I \circ \mathcal{L} = (-T^2) \circ I$ indicates that smoothness in \mathcal{L} translates into smoothness along $(-T^2)$.

Define $H_{T,+}^{1/2}(\partial_+ S\mathbb{M}) = \{w \in L_+^2, (-T^2)^{k/2} w \in L_+^2\}$ to obtain

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How far do these results generalize ?

The results above are sensitive to both the geometry and the boundary. In [Mishra-M. '19], [M. '19]: generalizations to geodesic disks of constant curvature, modeled over

$$M_{R,\kappa} = (\mathbb{D}_R, (1 + \kappa|z|^2)^{-2}|dz|^2), \quad R^2|\kappa| < 1.$$

Results: On $M_{R,\kappa}$, there is a weight function w such that $I_0: L^2(M_{R,\kappa}, w) \rightarrow L^2(\partial_+ SM_{R,\kappa})$ satisfies:

- $I_0^{-1} I_0$ is a C^∞ isomorphism.
- There are differential operators \mathcal{L} and $-T^2$ such that

$$I_0 w \circ \mathcal{L} = (-T^2) \circ I_0 w, \quad \text{and} \quad \mathcal{L}(I_0^{-1} I_0 w)^2 = id.$$

Open question: find more (simple !) surfaces where this works.

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Outline

- 1 Introduction
- 2 Main results
- 3 A numerical illustration
- 4 Elements of proof: mapping properties of unattenuated X-ray
 - The 'classical' functional setting $L^2(M) \rightarrow L^2_{\mu}(\partial_+ SM)$
 - The Euclidean disk: setting $L^2(\mathbb{D}) \rightarrow L^2(\partial_+ S\mathbb{D})$
- 5 Elements of proof: mapping properties of attenuated X-ray

Recalls

Let (M, g) be a simple Riemannian surface with geodesic vector field X , and $\Theta \in C^\infty(M, \mathfrak{u}(n))$ a 'Higgs field'. We define the attenuated X-ray transform $I_\Theta : L^2(M, \mathbb{C}^n) \rightarrow L^2_{\mu/\tau}(\partial_+ SM, \mathbb{C}^n)$ as

$$I_\Theta f = u|_{\partial_+ SM},$$

where $u : SM \rightarrow \mathbb{C}^n$ solves the transport equation

$$Xu + \Theta u = -f \quad (SM), \quad u|_{\partial_- SM} = 0.$$

Most recent results on the problem of recovering f from $I_\Theta f$ (case $n \geq 2$):

- Injectivity: [Paternain-Salo-Uhlmann '12].
- $L^2 - H^1$ stability estimate: [M.-Nickl-Paternain '20].

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Main theorems

Theorem (M'-Nickl-Paternain, 2020, arXiv preprint 2007.15892)

Let (M, g) a convex, non-trapping manifold with $\Theta \in C^\infty(M, \mathbb{C}^{n \times n})$. Then the operator $I_\Theta^* I_\Theta$ maps $C^\infty(M, \mathbb{C}^n)$ into itself.

Obtaining the converse mapping property currently requires strong assumptions on the background geometry + compact support.

Theorem (M'-Nickl-Paternain, 2020, arXiv preprint 2007.15892)

On M the Euclidean disk, let $\Theta \in C_c^\infty(M, \mathfrak{u}(n))$, and let the attenuated X-ray transform

$$I_\Theta: L^2(M, \mathbb{C}^n) \rightarrow L^2_{\mu/\tau}(\partial_+ SM, \mathbb{C}^n).$$

Then $I_\Theta^* I_\Theta$ is an isomorphism

$$C^\infty(M, \mathbb{C}^n) \xrightarrow{\cong} C^\infty(M, \mathbb{C}^n), \quad \tilde{H}^s(M, \mathbb{C}^n) \xrightarrow{\cong} \tilde{H}^{s+1}(M, \mathbb{C}^n), \quad s \geq 0.$$

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Elements of proof - forward mapping properties

Elements of proof - isomorphism properties

Conclusion

On the geodesic X-ray transform on the Euclidean disk (... and constant curvature disks)

- 1 Functional relations, link with degenerate elliptic operators.
- 2 Sharp mapping properties of I^*I and I , SVD of I for a special choice of weights on M and ∂_+SM .
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Perspectives:

- how far can we take 1-2 on simple surfaces ?
- higher dimensions ?
- case with non-trivial connections ?

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