

LECTURE 8: ENERGY METHODS

Friday, October 11, 2019

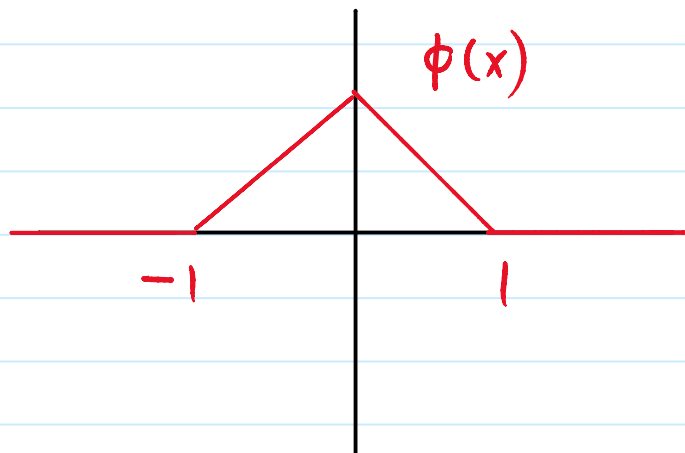
5:42 PM

I- THE PLUCKED STRING (section 2.1)

Example 2: [The plucked string]

$u_{tt} = u_{xx}$ ($c = 1$) with $u(x,0) = \phi(x)$ and $u_t(x,0) = 0$, where:

$$\phi(x) = \begin{cases} 0 & \text{for } x \leq -1 \\ 1 - |x| & \text{for } -1 < x < 1 \\ 0 & \text{for } x \geq 1 \end{cases}$$



D'Alembert says:

$$u(x,t) = 1/2 [\phi(x-t) + \phi(x+t)]$$

Now given the piecewise definition of ϕ , this becomes quite complicated, and we need to split this up into a lot of cases.

Let me illustrate the case $t = 1/2$:

$$u(x, 1/2) = 1/2 [\phi(x-1/2) + \phi(x+1/2)]$$

CASE 1: $x < -3/2$

Then $x - 1/2 < -2$ and $x + 1/2 < -1$

In that case $\phi(x-1/2) = 0$ and $\phi(x+1/2) = 0$, so

$$u(x, 1/2) = 1/2 (0 + 0) = 0$$

CASE 2: $-3/2 < x < -1/2$

Then $x - 1/2 < -1$ but $-1 < x + 1/2 < 0$

In that case $\phi(x-1/2) = 0$ but

$$\phi(x+1/2) = 1 - |x+1/2| = 1 - (-x-1/2) = 3/2 + x$$

$$\text{So } u(x, 1/2) = 1/2(0 + 3/2 + x) = 3/4 + x/2$$

CASE 3: $-1/2 < x < 1/2$

Then $-1 < x - 1/2 < 0$ but $0 < x + 1/2 < 1$

In that case:

$$\phi(x-1/2) = 1 - |x-1/2| = 1 + x - 1/2 = x + 1/2$$

$$\phi(x+1/2) = 1 - |x+1/2| = 1 - (x+1/2) = 1/2 - x$$

But then

$$u(x, 1/2) = 1/2 (x + 1/2 + 1/2 - x) = 1/2 (1) = 1/2$$

CASE 4: $1/2 < x < 3/2$

Then $0 < x - 1/2 < 1$ and $1 < x + 1/2 < 2$

$$\phi(x - 1/2) = 1 - |x - 1/2| = 1 - (x - 1/2) = 3/2 - x$$

$$\phi(x + 1/2) = 0$$

$$\text{So } u(x, 1/2) = 1/2 (3/2 - x) = 3/4 - x/2$$

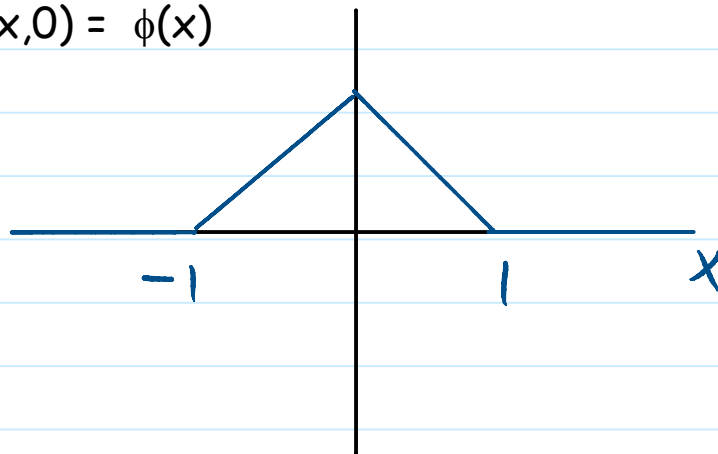
CASE 5: $x > 3/2$

Then $x - 1/2 > 1$ and $x + 1/2 > 2 > 1$, so

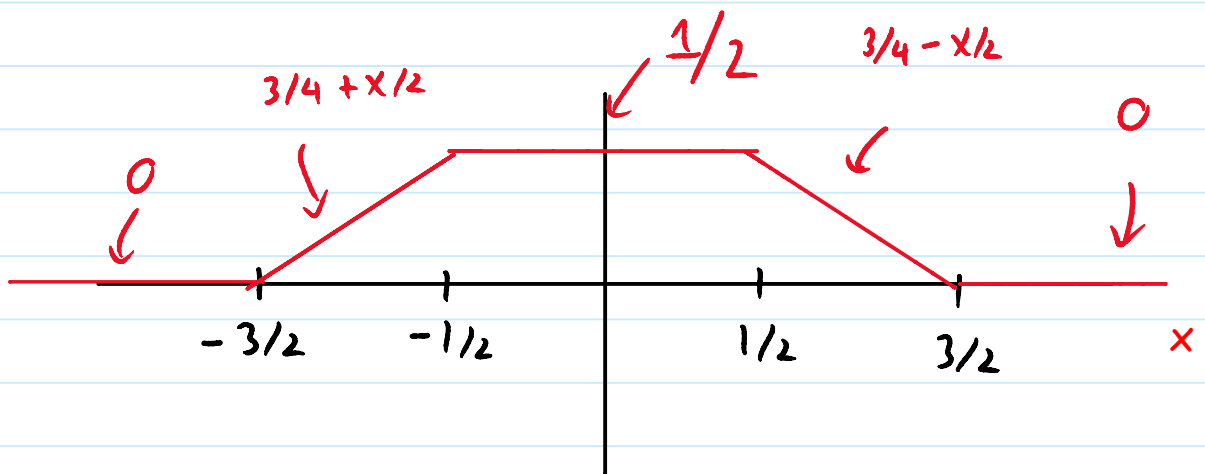
$$\phi(x - 1/2) = 0 \text{ and } \phi(x + 1/2) = 0, \text{ and so } u(x, 1/2) = 0$$

Picture: $u(x, t)$ for various t

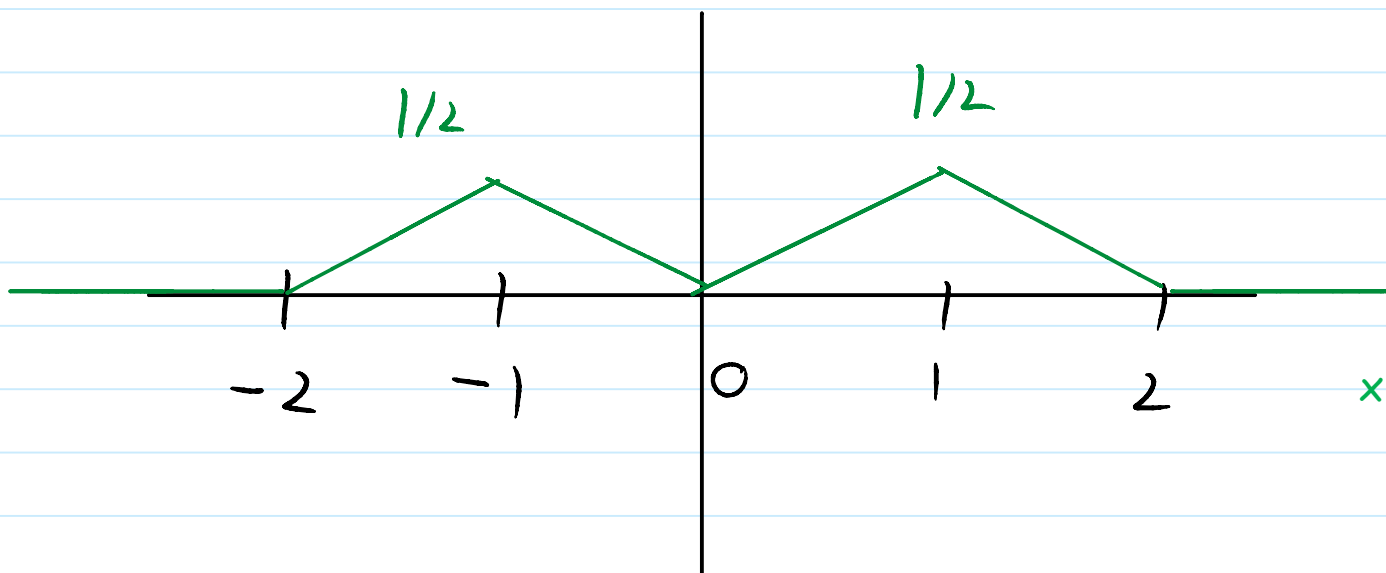
$$t = 0 \quad u(x, 0) = \phi(x)$$



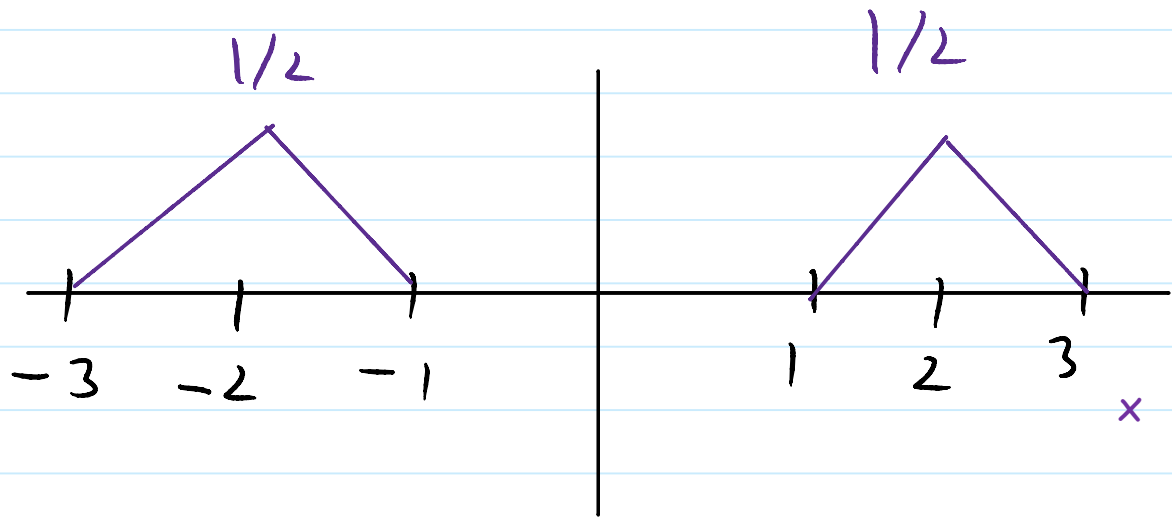
$t = 1/2$ Discussed above



$t = 1$ HW



$$t = 2$$



And so on and so forth!

This gives us a nice "movie" of the solution

II- THE ENERGY METHOD (Section 2.2)

Let's continue by discussing some more general properties of the wave equation.

Note: Everything below *any* solution of the wave equation. We are **NOT** using d'Alembert's formula here!

There are two main classes of PDE methods: Maximum Principle Methods (based on the maximum principle in 2.3) and Energy Methods (based on integration by parts).

Here, let me illustrate how the energy method works:

MAIN RESULT: [CONSERVATION OF ENERGY]

Suppose u solves $u_{tt} = u_{xx}$

Then the following energy $E(t)$ is conserved:

$$E(t) = \frac{1}{2} \int_{-\infty}^{\infty} (u_t)^2 + (u_x)^2 dx$$

Note: In physics, the first term is called the kinetic energy ($\frac{1}{2} m v^2$) and the second part the potential energy, so this says that the total energy is conserved.

Method 1: Show $E'(t) = 0$

Could go that route, and it's in fact easier, but it requires you beforehand to know what E is!

Method 2: Energy Method

Start with:

$$u_{tt} = u_{xx}$$

Multiply both sides by u_t

$$u_{tt} u_t = u_{xx} u_t$$

Now integrate with respect to x:

$$\underbrace{\int_{-\infty}^{\infty} u_{tt} u_t dx}_{(A)} = \underbrace{\int_{-\infty}^{\infty} u_{xx} u_t dx}_{(B)}$$

STUDY OF A:

Note that from calculus: $y'' y' = \frac{1}{2} [(y')^2]'$

$$\text{Hence: } u_{tt} u_t = \frac{1}{2} \frac{d}{dt} (u_t)^2$$

In particular:

$$\begin{aligned} A &= \int_{-\infty}^{\infty} \frac{1}{2} \frac{d}{dt} (u_t)^2 dx \\ &= \frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (u_t)^2 dx \end{aligned}$$

STUDY OF B:

Here we integrate by parts with respect to x

Note:

1) Integration by parts: $\int f'g = fg - \int fg'$

2) Here we assume the fg term is 0 (which basically means that our waves are 0 at $x = +/-$ infinity, which makes sense in practice)

3) $y y' = 1/2 (y^2)'$

$$\begin{aligned} \int_{-\infty}^{\infty} U_{xx} U_t dx &\stackrel{\text{IBP}}{=} - \int_{-\infty}^{\infty} U_x U_{xt} dx \\ &= - \int_{-\infty}^{\infty} \frac{1}{2} \frac{d}{dt} (U_x)^2 dx \\ &= - \frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (U_x)^2 dx \end{aligned}$$

$A = B$ then implies:

$$\frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (U_t)^2 = - \frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (U_x)^2 dx$$

$$\frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (U_t)^2 dx = -\frac{1}{2} \frac{d}{dt} \int_{-\infty}^{\infty} (U_x)^2 dx$$

$$\frac{d}{dt} \left[\underbrace{\frac{1}{2} \int_{-\infty}^{\infty} (U_t)^2 + (U_x)^2 dx}_{E(t)} \right] = 0$$

Therefore $E(t)$ is constant!

Remarks:

1) In particular $E(t) = E(0)$ and $E(0)$ only depends on our initial conditions ϕ and ψ

2) Application: Uniqueness of solutions

Suppose u and v both solve the wave equation

Let $w = u - v$, then w also solves the wave equation (check)
but with $\phi = 0$ and $\psi = 0$ (check)

Then, by # 1 in 2.2 (on HW), get $w = 0$, so $u - v = 0$ so $u = v$