

# LECTURE 9: THE HEAT EQUATION (I)

Monday, October 14, 2019 5:23 PM

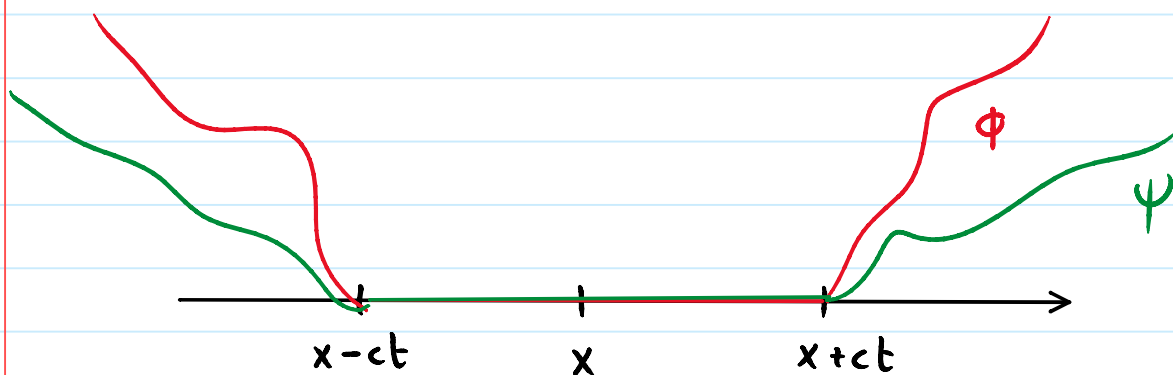
## I- DOMAIN OF DEPENDENCE (section 2.2)

Let's look again at D'Alembert's formula:

$$u(x,t) = \frac{1}{2} (\phi(x-ct) + \phi(x+ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds$$

Notice that the value of  $u$  at  $(x,t)$  only depends on the values of  $\phi$  and  $\psi$  on the interval  $[x-ct, x+ct]$ .

In particular, suppose  $\phi = 0$  and  $\psi = 0$  on the interval  $[x-ct, x+ct]$  but are nonzero elsewhere, as in this picture:

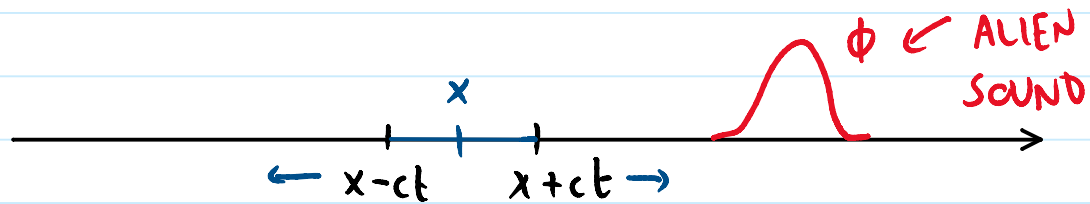


Then  $u(x,t) = 0$

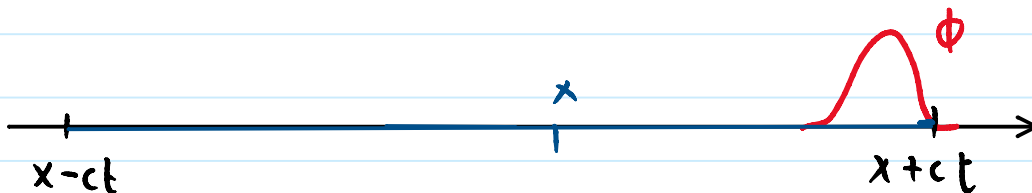
In other words, anything **outside** the interval  $[x-ct, x+ct]$  doesn't affect  $u$  at all!

**Interpretation:** The wave equation has **finite** speed of propagation: If an alien far far away (outside the interval  $[x-ct, x+ct]$ ) yells at you, you won't really hear the sound until time is so large that  $\phi$  or  $\psi$  is nonzero on  $[x-ct, x+ct]$ .

Ex:  $t$  small: Here  $u(x,t) = 0$



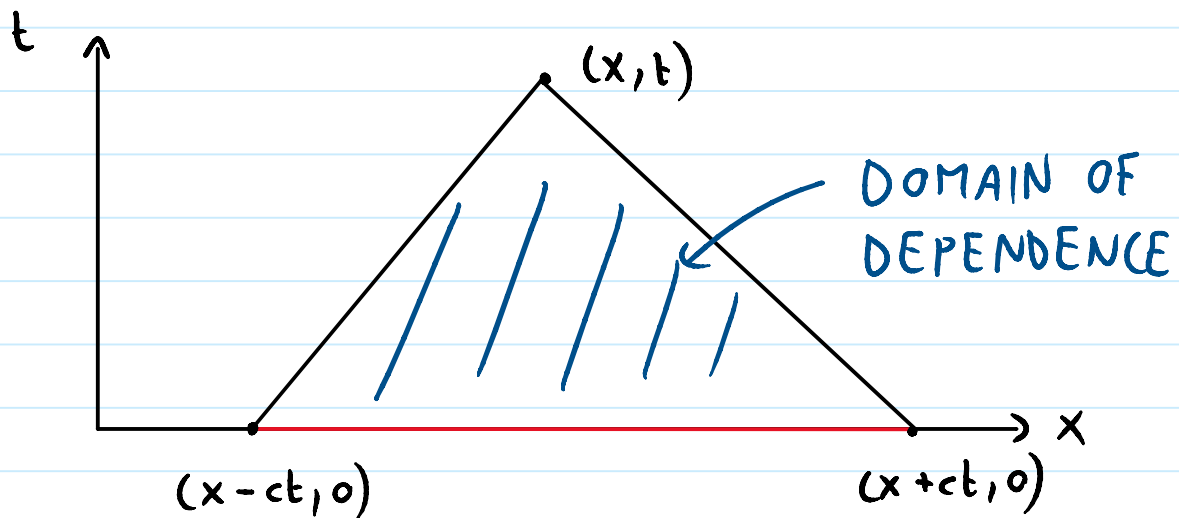
$t$  large,  $u(x,t)$  is nonzero:



We can represent this pictorially as follows:

Start with  $(x,t)$  and connect it with the points  $(x-ct, 0)$  and  $(x+ct, 0)$  (Why?  $\phi(x-ct) = u(x-ct, 0)$ )

Picture:



You get a triangle called the "domain of dependence"

**Interpretation:** If  $\phi$  and  $\psi$  are both 0 on the bottom leg of the triangle, then  $u = 0$  at the top vertex (and even on the whole triangle!)

**Note:** In higher dimensions, the triangle becomes a cone.

This ends our 'first look' at the wave equation! We'll come back to it very soon, but now let's discuss our second important PDE: The Heat Equation!

## II- THE HEAT EQUATION

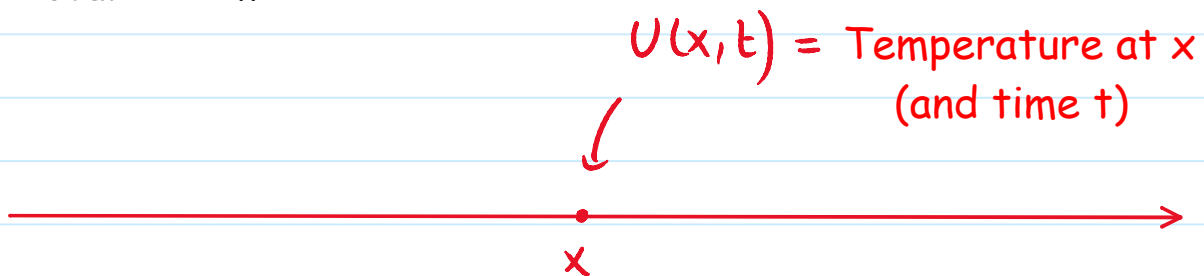
$$u = u(x,t)$$

$$u_t = k u_{xx}$$

$$u(x,0) = \phi(x)$$

Also known as the diffusion equation,  $u(x,t)$  measures the temperature of an (infinite) metal rod at position  $x$  and time  $t$

**Picture:** Time =  $t$



## III- FUNDAMENTAL SOLUTION (Section 2.4)

Let's first focus on finding a solution (section 2.4), then then move on to discussing general properties (section 2.3).

**WARNING:** Our approach is VERY different from the wave equation. Here, instead of finding the general solution, we will first find **one** solution, and then build the general solution from that.

**Goal:** Find a solution of  $u_t = u_{xx}$  ( $k = 1$ )

In this case, we will look for a solution of a special form

## STEP 0

Silly trick: If you ignore  $u$  in  $u_t = u_{xx}$  you get:

$$t = x^2$$

$$\Rightarrow \frac{x^2}{t} = 1$$

$$\Rightarrow \left( \frac{x}{\sqrt{t}} \right)^2 = 1$$

So it might seem like the quantity  $y = \frac{x}{\sqrt{t}}$

might be useful, and indeed it is!

## STEP 1

Guess:  $u(x,t) = v\left(\frac{x}{\sqrt{t}}\right)$

For some function  $v = v(y)$  to be found

This doesn't *quite* work, unfortunately.

Better guess:

$$u(x,t) = \frac{1}{t^\alpha} v\left(\frac{x}{\sqrt{t}}\right)$$

$$u(x,t) = t^{-\alpha} v\left(x t^{-1/2}\right)$$

For some  $\alpha$  and  $v$  to be determined.

(Intuitively:  $u$  "blows up" near  $t = 0$ )

## STEP 2

Plug our guess in our PDE  $u_t = u_{xx}$

$$\begin{aligned} u_x &= \left( t^{-\alpha} v\left(x t^{-1/2}\right) \right)_x \\ &= t^{-\alpha} v'\left(x t^{-1/2}\right) \left(t^{-1/2}\right) \\ &= t^{-\alpha - \frac{1}{2}} v'\left(x t^{-1/2}\right) \end{aligned}$$

$$\begin{aligned} u_{xx} &= t^{-\alpha - \frac{1}{2}} v''\left(x t^{-1/2}\right) \left(t^{-1/2}\right) \\ &= t^{-\alpha - 1} v''(\gamma) \end{aligned}$$

$$\begin{aligned} u_t &= \left( t^{-\alpha} v\left(x t^{-1/2}\right) \right)_t \\ &= -\alpha t^{-\alpha - 1} v\left(\frac{x}{\sqrt{t}}\right) + t^{-\alpha} v'\left(\frac{x}{\sqrt{t}}\right) \left(x t^{-1/2}\right)_t \end{aligned}$$

$$= -\alpha t^{-\alpha-1} V(\gamma) + t^{-\alpha} V'(\gamma) \left(-\frac{x}{2} t^{-3/2}\right)$$

$$= -\alpha t^{-\alpha-1} V(\gamma) - \frac{x}{2} t^{-\alpha-\frac{3}{2}} V'(\gamma)$$

So  $u_t = u_{xx}$  implies:

$$-\alpha \cancel{t^{-\alpha-1}} V - \frac{x}{2} \cancel{t^{-\alpha-\frac{3}{2}}} V' = \cancel{t^{-\alpha-1}} V''$$

$$-\frac{\alpha}{\cancel{t}} V - \frac{x}{2 \cancel{t} \sqrt{t}} V' = \frac{1}{\cancel{t}} V''$$

$$-\alpha V - \underbrace{\left(\frac{x}{\sqrt{t}}\right)}_{\frac{\gamma}{2}} V' = V''$$

$$\Rightarrow V'' + \frac{\gamma}{2} V' + \alpha V = 0$$

**Note:** If we didn't include the  $1/t^\alpha$  term, we wouldn't have such a nice simplification with the  $\gamma$  (and instead have weird  $x$  and  $t$  terms)

At this point we are stuck, but remember that we still have a choice in our constant  $\alpha$ !

### STEP 3

Choose  $\alpha = 1/2$  (this allows us to factor out the second and third term)

Then we get:

$$V'' + \frac{\gamma}{2} V' + \frac{1}{2} V = 0$$

$$V'' + \frac{1}{2} (\gamma V' + V) = 0$$

$$V'' + \frac{1}{2} (\gamma V)' = 0$$

$$[V' + \frac{1}{2} \gamma V]' = 0$$

$$\Rightarrow V' + \frac{\gamma}{2} V = C$$

Assume  $C = 0$

(Why? We just need to find **ONE** solution. Moreover, in physical models,  $v(y)$  and  $v'(y)$  goes to 0 as  $y$  goes to infinity)

$$\Rightarrow V' + \frac{\gamma}{2} V = 0$$



STEP 4 Solve this!

$$v' = -\frac{\gamma}{2} v$$

$$\Rightarrow \frac{v'}{v} = -\frac{\gamma}{2}$$

$$\Rightarrow [\ln |v|]' = -\frac{\gamma}{2}$$

$$\Rightarrow \ln |v| = -\frac{\gamma^2}{4} + C$$

$$\Rightarrow |v| = e^{-\frac{\gamma^2}{4} + C}$$

$$\Rightarrow v = \underbrace{\pm e^C}_{\text{Arbitrary}} e^{-\frac{\gamma^2}{4}}$$

$$\Rightarrow v(\gamma) = C e^{-\frac{\gamma^2}{4}}$$

STEP 5: CONCLUSION

$$\begin{aligned} u(x,t) &= \frac{1}{t^{1/2}} v\left(\frac{x}{\sqrt{t}}\right) \quad (\alpha = \frac{1}{2}) \\ &= \frac{1}{\sqrt{t}} C e^{-\frac{1}{4} \left(\frac{x}{\sqrt{t}}\right)^2} \end{aligned}$$

$$u(x,t) = \frac{C}{\sqrt{t}} e^{-\frac{x^2}{4t}}$$

Finally, choose  $C = \frac{1}{\sqrt{4\pi}}$  (see below why), then

$$\frac{C}{\sqrt{t}} = \frac{\frac{1}{\sqrt{4\pi}}}{\sqrt{t}} = \frac{1}{\sqrt{4\pi t}}$$

**Definition:** The **fundamental solution**  $S(x,t)$  of the heat equation  $u_t = u_{xx}$  is:

$$S(x,t) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}$$

**Remarks:**

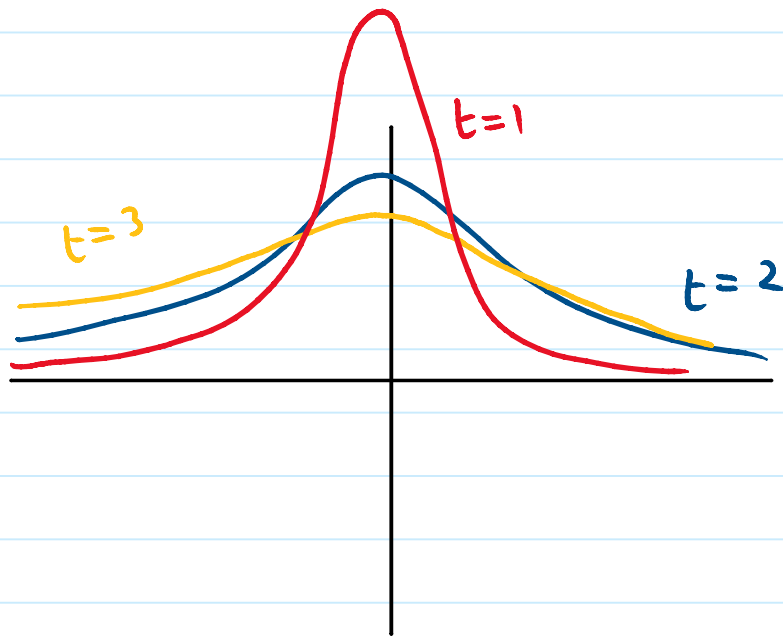
- 1) The reason we chose  $C$  as we did is because, with this choice of  $C$ , we have:

$$\int_{-\infty}^{\infty} S(x,t) dx = 1$$

That is, for each  $t$ , the area under the solution is 1 (see # 6 and 7 in 2.4 on HW # 4)

- 2)  $S(x,t)$  looks like the bell curve  $\exp(-x^2)$ , but more and more spread out as  $t$  gets larger. At each time, the area under the curve is 1.

Picture:



- 3) To solve  $u_t = k u_{xx}$  replace  $t$  by  $kt$  to get that the fundamental solution is:

$$S(x,t) = \frac{1}{\sqrt{4\pi kt}} e^{-\frac{x^2}{4kt}}$$

This is the version we'll be dealing with throughout the course.