LECTURE 18: SEPARATION OF VARIABLES (I)

Monday, November 4, 2019 12:47 PM

Welcome to the most important PDE technique of the course: Separation of variables! It's **the** technique that will help us solve all the remaining PDEs (like the rod of finite length, and even Laplace's equation)

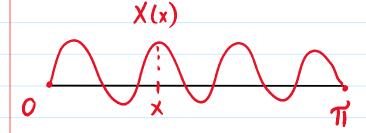
LONG technique, but will do this at least 4 times!

I- BOUNDARY VALUE PROBLEM

Before we do that, let's solve a problem that is purely ODE, but will be *crucial* in our process.

Example: Find a nonzero solution X(x) of

$$\begin{cases} X''(x) = \lambda X(x) \\ X(0) = 0 \\ X(\pi) = 0 \end{cases}$$



Idea: The solutions of, for example, X'' = -X (which involve cos & sin) are very different from the solutions of X'' = X (which involve exponential terms), so we need to argue in terms of the sign of λ

CASE 1: $\lambda > 0$

Then $\lambda = \omega^2$ for some $\omega > 0$

(Ex: If λ = 4, then λ = 4 = 2² so ω = 2)

Why: If you don't assume that, you'll get nasty $\sqrt{\lambda}$ everywhere

Then:

$$X'' = \lambda X \Rightarrow X'' = \omega^2 X \Rightarrow X'' - \omega^2 X = 0$$

From Math 3D (ODE), we get:

Auxiliary equation: $r^2 - \omega^2 = 0 \Rightarrow r^2 = \omega^2 \Rightarrow r = \pm \omega$

But
$$X(0) = A e^{\omega 0} + B e^{-\omega 0} = A + B = 0$$
 (since $X(0) = 0$)

So
$$X(x) = A e^{\omega x} - A e^{-\omega x}$$

But
$$X(\pi) = 0 \Rightarrow A e^{\omega \pi} - A e^{-\omega \pi} = 0$$

 $\Rightarrow A(e^{\omega \pi} - e^{-\omega \pi}) = 0$
 $\Rightarrow e^{\omega \pi} - e^{-\omega \pi} = 0$

(A \neq 0, because otherwise X(x) = A $e^{\omega x}$ - A $e^{-\omega x}$ = 0)

$$\Rightarrow$$
 $e^{\omega \pi} = e^{-\omega \pi}$

$$\Rightarrow \omega \pi = -\omega \pi$$

$$\Rightarrow \omega = 0$$

But then $\lambda = \omega^2 = 0$ which is not positive! =×=

CASE 2: $\lambda = 0$

Then X'' = 0 X => X''(x) = 0
=>
$$X(x) = Ax + B$$

$$X(0) = A0 + B = B = 0$$
, so $X(x) = Ax$

$$X(\pi) = A\pi = 0 \Rightarrow A = 0$$
, but then $X(x) = 0x = 0 \Rightarrow 0$

CASE 3: λ < 0

Then $\lambda = -\omega^2$ for some w > 0

(Ex: $\lambda = -4 = -2^2$)

$$X'' = \lambda X \Rightarrow X'' = -\omega^2 X \Rightarrow X'' + \omega^2 X = 0$$

Aux:
$$r^2 + \omega^2 = 0 \Rightarrow r^2 = -\omega^2 \Rightarrow r = \pm \omega i$$

$$X(x) = A \cos(\omega x) + B \sin(\omega x)$$

So
$$X(x) = 0 \cos(\omega x) + B \sin(\omega x) = B \sin(\omega x)$$

$$X(\pi) = S \sin(\omega \pi) = 0$$

=> $\sin(\omega \pi) = 0$
=> $\omega \pi = \pi m \ (m = 1, 2, ...)$
=> $\omega = m$

Answer: For every m, we have a solution,

$$X(x) = \sin(\omega x) = \sin(mx) \quad (m = 1, 2, ...)$$

Note: This problem is sometimes called an eigenvalue problem.

The eigenvalues of X'' = λ X are λ = $-\omega^2$ = $-m^2$ And the eigenfunctions are X(x) = $\sin(mx)$

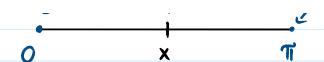
(Compare to $Ax = \lambda x$, from linear algebra, here have infinitely many eigenvalues and eigenvectors)

II- SETTING

Now let's go back to PDEs and see why we need the problem above

Example: Solve

$$\begin{cases} u_{t} = k u_{xx} & (0 < x < \pi) \\ u(0,t) = 0, u(\pi,t) = 0 \\ u(x,0) = x^{2} \end{cases}$$



It's the usual finite rod problem that's insulated at the endpoints, but this time we actually want to find a solution.

Remarks:

- We already showed uniqueness using energy methods/maximum principle, so everything just boils down to finding A solution.
- 2) We want to find NONZERO solutions (the zero solution isn't very interesting)

Motivation: For the heat equation on the whole line, we assumed u had a special form

$$u(x,t) = 1/\sqrt{t} v(x/\sqrt{t})$$

And we then turned the heat equation into an ODE for v:

$$v'' + (y/2) v' + (1/2) v = 0 \leftarrow MUUUCH simpler!$$

Here we want to do something similar.

III- SEPARATION OF VARIABLES

STEP 1: Separation of variables

1) Assume u(x,t) has a special form, namely:

$$u(x,t) = X(x) T(t)$$
 (*)

Where X is a function of x and T is a function of t only

Ex: sin(x) cos(t) is of this form, and so is $(x^2 + 1)(t^3 - 4)$, but ln(3x + 2t) is not of this form!

2) Plug (*) into $u_t = k u_{xx}$

$$(X(x) T(t))_t = k (X(x) T(t))_{xx}$$

$$X(x) (T(t))_t = k (X(x))_{xx} T(t)$$

$$X(x) T'(t) = k X''(x) T(t)$$

3) Put all the X terms on one side and all the T terms on the other side

$$\frac{T'(t)}{k} = \frac{X''(x)}{X(x)}$$

Always put any extra terms like k (or c²) with T
(because we want to keep the X equation as simple as possible)

STEP 2: Constant

1) Look at the quantity
$$\frac{X''(x)}{X(x)} = \frac{T'(t)}{k}$$

 $X(x)$

On the one hand,

$$\left(\frac{X''(x)}{X(x)}\right)_{t} = 0$$

But also

$$\left(\frac{X''(x)}{X(x)}\right)_{x} = \left(\frac{T'(t)}{kT(t)}\right)_{x} = 0$$

So the $\frac{X''(x)}{X(x)}$ doesn't depend on x or t, so it is constant!

Conclusion:
$$\frac{X''(x)}{X(x)} = \frac{T'(t)}{kT(t)} = constant = \lambda$$

$$\Rightarrow \underline{X''(x)} = \lambda \Rightarrow X''(x) = \lambda X(x)$$

$$X(x)$$

And
$$\underline{T'(t)} = \lambda \Rightarrow T'(t) = k\lambda T(t)$$

kT(t)

GOOD! We're on the right track; instead of a PDE, we get 2 ODEs.

2) Focus on X(x) equation! (don't touch T until the end!!!)

So far:
$$X''(x) = \lambda X(x)$$

Now use the boundary conditions:

$$u(0,t) = 0 \Rightarrow X(0) T(t) = 0 \text{ (since } u(x,t) = X(x)T(t))$$

=> $X(0) = 0$

(Why can you cancel out T? If T = 0, then u = XT = 0, which wouldn't be interesting. So pick a t where T(t) is nonzero, and cancel out T there)

Similarly
$$u(\pi,t) = 0 \Rightarrow X(\pi)T(t) = 0 \Rightarrow X(\pi) = 0$$

Hence we get the ODE

$$X''(x) = \lambda X(x)$$

 $X(0) = 0$
 $X(\pi) = 0$

Same ODE as the beginning of lecture!!!

STEP 3: Boundary-value problem

INSERT PREVIOUS EXAMPLE

Conclusion: $\lambda = -m^2$ (m = 1, 2, ...)

$$X(x) = \sin(mx)$$
 (m = 1, 2, ...)

STEP 4: T equation

$$\frac{T'}{kT} = \lambda = -m^2$$

$$=> T' = -m^2 k T$$

=>
$$T(t) = e^{-m^2 kt}$$
 (ignore constant C for now)

Conclusion: For every m = 1, 2, ...

$$u(x,t) = X(x)T(t) = \sin(mx) e^{-m kt}$$
 is a solution of our PDE

STEP 5: Linearity

Important observation: Since the PDE is linear, any linear combo of the above solution is also a solution!

So

$$u(x,t) = A_1 \sin(1x) e^{-(1)kt} + A_2 \sin(2x) e^{-(2)kt} + ...$$

$$u(x,t) = \sum_{M=1}^{\infty} A_{m} \sin(mx) e^{-m^{2}kt}$$
 solves the PDE

STEP 6: Initial Condition

$$u(x,0) = \sum_{M=1}^{2} A_{m} \sin(mx) e^{-m k0} = \sum_{M=1}^{2} A_{m} \sin(mx)$$

$$x^{2} = \sum_{M=1}^{\infty} A_{m} \sin(mx)$$

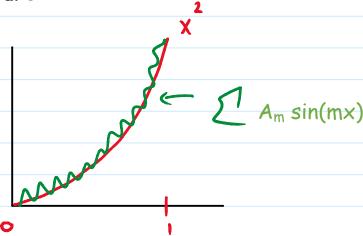
=> BIG QUESTION:

Can you find A_m such that the above is true? In other words, can you write x^2 as a linear combo of sines?

YES, see Chapter 5

(In fact, this is precisely why Fourier series were invented)





IV- EASIER PROBLEM

$$u_{t} = k u_{xx}$$

$$u(0,t) = 0, u(\pi,t) = 0$$

$$u(x,0) = 3 \sin(2x) + 4 \sin(3x)$$

Steps 1-5 still work:

$$u(x,t) = \int_{M=1}^{\infty} A_m \sin(mx) e^{-m kt}$$

This time:

$$u(x,0) = \sum_{M=1}^{\infty} A_{m} \sin(mx) = 3 \sin(2x) + 4 \sin(3x)$$

=>
$$A_1 \sin(x) + A_2 \sin(2x) + A_3 \sin(3x) + ...$$

= $0 \sin(x) + 3 \sin(2x) + 4 \sin(3x)$

$$A_1 = 0$$
, $A_2 = 3$, $A_3 = 4$, $A_4 = 0$, etc.

(just compare the coefficients)

$$u(x,t) = \int_{M=1}^{2} A_m \sin(mx) e^{-m^2 k t}$$

=
$$3 \sin(2x) e^{-2kt} + 4 \sin(3x) e^{-3kt}$$

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$$u(x,t) = 3 \sin(2x) e^{-4kt} + 4 \sin(3x) e^{-9kt}$$

Interpretation: Solution starts as $3 \sin(2x) + 4 \sin(3x)$, but eventually dies down because the exponential terms go to 0



